# The Structure of Overcrowded Aromatic Compounds. Part III.* The Crystal Structure of 3:4-Benzophenanthrene. 

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The crystal structure of $3: 4$-benzophenanthrene has been analysed by two-dimensional Fourier methods. The molecular orientation and position in the unit cell were fixed by a combination of the molecular Fourier transform method and the interpretation of non-space-group absences and near absences of certain zonal reflections. The atomic co-ordinates were refined by successive $\mathrm{F}_{\mathrm{o}}-\mathrm{F}_{\mathrm{c}}$ syntheses to $R$ factors of $0.22,0.23$, and $\mathbf{0 . 2 6}$ in the $h k 0,0 k l$, and $h 0 l$ zones.

The final co-ordinates can be fitted, within a root-mean-square deviation of $0.03 \AA$, to a molecular model of symmetry 2 , in which a separation of $3.0 \AA$ between the centres of the overcrowded carbon atoms labelled 1 and 12 in Fig. 1 is achieved by distortion of the benzene rings. This model consists of two pairs of six-membered rings, related by a twofold axis, which are folded about their long diagonals through angles of $10^{\circ}$ and $14^{\circ}$.

The distortions of the molecular skeleton have been related to the difference between the calculated and measured resonance energies and to the non-vanishing dipole moment; the nearly tetrahedral bond angles at four of the secondary carbon atoms to the close approach distance of $3 \cdot 13 \AA$ between two such atoms of neighbouring molecules, and to the preferential attack on these same atoms by electrophilic reagents.
The analysis of the structure of $3: 4$-benzophenanthrene was undertaken in continuation of work outlined in Part I ( $J$., 1954, 3288). The distance of $2 \cdot 4 \AA$ between carbon atoms labelled 1 and 12 in Fig. $1 \dagger$ in a conventional model of the molecule is shorter than the approach distances found in any of the overcrowded molecules discussed in Part I, and one may therefore expect the benzophenanthrene ring system to be distorted in order to achieve relief from overcrowding. To investigate quantitatively any such deformation we have determined the structure of $3: 4$-benzophenanthrene by $X$-ray crystallography, and now report the results of an analysis by two-dimensional Fourier methods.

A preliminary crystallographic investigation has been reported by Iball ( $Z$. Krist., $1938,99, A, 230$ ), from which we cite the following data : $a=14 \cdot 60, b=14 \cdot 00, c=5 \cdot 76 \AA$;

[^0]$d$ (found) $1 \cdot 265, d$ (calc.) $1 \cdot 278 ; n=4$. Absences : $h 0 l$ absent for $h+l$ odd, $0 k 0$ absent for $k$ odd. Possible space-groups : $P 2_{1} 2_{1} 2_{1}, P m n, P m n m$, with $P m n$ favoured on evidence referred to below.

The cell dimensions and space-group were now determined from zero- and first-layer Weissenberg photographs about the three principal axes. The lattice constants were corrected by extrapolation to $\theta=90^{\circ}$ (Buerger, " $X$-Ray Crystallography," J. Wiley and Sons, New York, 1942, Ch. 20) and are accurate to $0 \cdot 1 \%$ : $a=14 \cdot 69, b=14 \cdot 19, c=$ $5 \cdot 76 \AA ; d$ (calc.) $=1 \cdot 261$. Absences : $h 00$ absent for $h$ odd, $0 k 0$ absent for $k$ odd, $00 l$ absent for $l$ odd. Space-group : $P 2_{1} 2_{1} 2_{1}$.

The two sets of lattice parameters are in good agreement. Our Weissenberg photographs show that the $h 0 l$ reflections with $h+l$ odd are weak but not systematically absent. Thus the extinctions referred to by Iball as systematic are in fact pseudo-absences; the space-group is therefore unequivocally determined as $P 2_{1} 2_{1} 2_{1}$. The significance of these pseudo-extinctions for the molecular shape and arrangement is discussed in detail below.

Fig. 1. The planar model of the 3:4benzophenanthrene molecule, showing the molecular axes and the nomenclature of atoms and rings used in this paper.


Fig. 2. The (010) projections of the space-groups (a) $P 2_{1} 2_{1} 2_{1}$ and (b) Pmn2 $2_{1}$. Possible molecular arrangements in.$P 2_{1} 2_{1} 2_{1}$ with the additional symmetry implied by the absences of hol spectra with $h+l$ odd are shown in c and d .

In common with other workers, we prefer, for space-group $P 2_{1} 2_{1} 2_{1}$, to choose origins for the three principal projections at the points of intersection of the corresponding twofold screw axes with the respective planes of projection. Hence, the origins of the (001), (100), and ( 010 ) projections are at the points $\left(\frac{1}{4} 00\right)$, ( $0 \frac{1}{4} 0$ ), and ( $00 \frac{1}{4}$ ), respectively, in the coordinate system of the " Internationale Tabellen zur Bestimmung von Kristallstrukturen" (Gebr. Borntraeger, Berlin, 1935). For the sake of uniformity all co-ordinates mentioned in the following discussion are referred to the origin ( $\frac{1}{4} 00$ ).

The length of the $c$ axis permits only a single molecule to be accommodated along [001], with the molecular plane inclined at about $45^{\circ}$ to (001). For additional information as to the molecular orientation, packing considerations are not very helpful because of the similarity of the $a$ and $b$ axes and the relatively symmetrical shape of the molecule; nor is the level of the unitary structure factors sufficiently high to favour the use of any of the analytical methods of sign determination. The correct molecular orientation was eventually derived by arguments based mainly on the molecular Fourier transform method and on the interpretation of the pseudo-absences in the zonal reflections.

In applying the Fourier transform method we have followed the procedure outlined by Klug (Acta Cryst., 1950, 3, 176) in the analogous instance of triphenylene and shall therefore restrict the following discussion to features peculiar to the present problem. The real and imaginary parts of the Fourier transform of 3:4-benzophenanthrene were calculated on the
assumption of a planar molecule with atoms at the corners of regular hexagons of side $1.40 \AA$. The system of molecular axes is shown in Fig. 1.

Comparison of the weighted reciprocal lattice with the Fourier transform led to the conclusions that (1) $a^{*}$ must lie along the $O L$ axis of the transform, (2) $b^{*}$ must lie at about $45^{\circ}$ to the plane of the transform, (3) the approximate co-ordinates ( $u_{0}, v_{0}$ ) of the molecular reference point $O$ are either $(0.01,0.25)$ or $(0.26,0.25)$, (4) the poor agreement between the Fourier transform and observed values of the structure factors beyond the low orders of $\mathrm{F}(h k 0)$ suggests that the planar and regular model used for the computation of the transform represents only inadequately the structure of the real molecule.

A second clue to the molecular orientation is provided by the pseudo-absences in the $h 0 l$ zone, which imply that the $b$ axis projection of the unit cell approximates to that of space-group $\mathrm{Pmn2}_{1}$ (Figs. 2, $a$ and $b$ ). There are two possible molecular arrangements in $P 2_{1} 2_{1} \mathbf{2}_{1}$ satisfying this condition of additional symmetry, of which one can be decisively


Fig. 3. The electron-density projection on (001), calculated with 144 out of 154 observed terms. The contours are drawn at intervals of $1 \mathrm{e} . \AA^{-2}$, with the one-electron line dotted.
ruled out. This orientation (Fig. 2, c) requires the projected molecule to have a pseudocentre of symmetry at the point $u=0, w=0$. As the only possible centrosymmetric projection of benzophenanthrene is that down the molecular $O L$ axis it follows that $O L$ would necessarily lie parallel to $b$. This arrangement is contradicted by the evidence of the Fourier transform, according to which $O L$ is parallel to $a$. Moreover, it is ruled out by the value of $F(230)$, which it would require to be small whereas in fact this reflection is the second strongest in the ( $h k 0$ ) zone. In his discussion of possible space-groups Iball put forward this molecular arrangement in $P 2_{1} 2_{1} 2_{1}$ as the explanation for the extra $h 0 l$ absences; he rejected it as an unlikely packing arrangement and concluded that the spacegroup could not be $P 2_{1} 2_{1} 2_{1}$.

The second possible molecular arrangement (Fig. 2, d) requires the projected molecule to have a pseudo-mirror plane parallel to (001) at $w=0.25$. This must coincide with the projected $O L$ axis, evidently a molecular twofold axis of symmetry. Thus, the information deduced from the pseudo-absences, together with the evidence from the Fourier transform, leads to the following conclusions: (1) the molecule has an approximate twofold axis of symmetry along $O L$, nearly parallel to the $a$ axis; (2) $O M$ is inclinded at about $45^{\circ}$ to the $a$ axis; (3) the approximate co-ordinates of the point $O$ are $u_{0}=0.01$ or 0.26 (Fourier transform), $v_{0}=0.25$ (Fourier transform), $w_{0}=0.25$ ( $h 0 l$ pseudo-absences).

Conditions (1) and (3) together imply that the symmetry of the cell approximates to that of space-group $B 22_{1} 2$; we should thus expect the pseudo-absences to extend throughout the reciprocal lattice. Indeed, Weissenberg photographs of the $h k 1$ and $h k 4$ spectra as well as of the three principal zones reveal that reflections with $h+l$ odd tend to be generally weaker than those with $h+l$ even.

The ambiguity in the $u_{0}$ parameter was resolved in favour of $u_{0}=0.01$ from considerations of molecular packing; during the course of the analysis the second alternative was tested and ruled out on the basis of the $h 0 l$ structure factors.

Atomic co-ordinates were calculated for a planar and regular molecule with $u_{0}=0.01$, $v_{0}=0.25, w_{0}=0.25, O L$ parallel to [100], and $O M$ at $45^{\circ}$ to ( 001 ). Refinement from this original set of co-ordinates was carried out by means of successive $F_{O}-F_{C}$ syntheses until the $R$ factor had decreased to $0 \cdot 22$. There remained at this stage ten $h k 0$ structure factors

Fig. 4. The electron-density projection on (100).

of uncertain sign, of which six had unitary scattering factors exceeding 0.05 , the largest being $U(\mathbf{1 4}, 2,0)=\mathbf{0} \cdot 16$. However, as a further difference synthesis failed to improve the agreement for these reflections it was concluded that the atomic co-ordinates already determined were the best that could be derived from the present experimental data.

A Fourier summation of 144 terms, out of 154 observed, was calculated at intervals of 120ths along $a$ and $b$ (Fig. 3). This projection shows that the molecule is not regular and probably not planar. By means of an adjustable wire model a non-planar molecular configuration was fitted to this projection, in which the mean plane of the molecule was inclined at about $45^{\circ}$ to the projection plane. With the molecular reference point $O$ located at $w_{0}=0 \cdot 25$, approximate $w$ co-ordinates were obtained from this model, giving $R$ factors for the $0 k l$ and $h 0 l$ zones of 0.42 and 0.34 , respectively. Although the degree of overlap is smaller in the $0 k l$ than in the $h 0 l$ zone, the latter was refined first because of the smaller initial value of its $R$ factor. After five difference syntheses $R(h 0 l)$ had been reduced to $0 \cdot 18$; however, $R(0 k l)$ had risen to 0.47 and this set of $w$ co-ordinates was therefore discarded.

Adjustment of the initial set of $w$ co-ordinates produced a new molecular model which
gave values of $R(0 k l)$ and $R(h 0 l)$ of 0.26 and 0.30 for reflections up to $2 \sin \theta=1 \cdot 0$. Refinement of this model by difference syntheses on the $0 k l$ reflections resulted after eight cycles in values for $R(0 k l)$ and $R(h 0 l)$ of 0.023 and 0.26 , at which point refinement ceased. The final Fourier projection on to (100), embodying 58 of the 62 observed 0 kl structure factors, is shown in Fig. 4, and Table 1 lists the final $z$ co-ordinates in col. 4 and the spurious $z$ co-ordinates, derived from the $h 0 l$ zone only, in col. 5 .

No attempt has been made to improve the rather high discrepancy factors ( $0 \cdot 22,0 \cdot 23$, and 0.26 in the $h k 0,0 k l$, and $h 0 l$ zones), which may be partly due to the exclusion of the hydrogen contributions from the calculated structure factors. That these contributions are not insignificant is shown by the appearance of a (001) projection obtained by one of us on $X$-RAC; this summation contains all observed $h k 0$ structure factors, including the ten terms referred to above, whose signs were determined by means of the " negative-background criterion " (e.g., Pepinsky, " Computing Methods and the Phase Problem in X-Ray Crystal Analysis," State College. Pa., 1952, p. 3). Although the carbon peaks are not


Fig. 5. The projection of the unit-cell contents on (001).

Table 1. Experimentally-determined atomic co-ordinates (in $\AA$ ).

| Atom* | $x$ | $y$ | $z$ | Spurious $z$ co-ordinates | Atom * | $x$ | $y$ | $z$ | Spurious $z$ co-ordinates |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 1 | $1 \cdot 81$ | $2 \cdot 06$ | $0 \cdot 86$ | 0.94 | 10 | $1 \cdot 60$ | $5 \cdot 97$ | $4 \cdot 46$ | $4 \cdot 06$ |
| 2 | $2 \cdot 39$ | $1 \cdot 18$ | 0.05 | -0.02 | 11 | $2 \cdot 19$ | $5 \cdot 43$ | $3 \cdot 34$ | $2 \cdot 86$ |
| 3 | $1 \cdot 70$ | 0.65 | $-1.25$ | $-1 \cdot 10$ | 12 | 1.73 | $4 \cdot 64$ | $2 \cdot 40$ | 2.00 |
| 4 | $0 \cdot 50$ | 1.08 | $-1.38$ | $-1.38$ | 13 | $0 \cdot 40$ | $2 \cdot 54$ | 0.58 | 0.71 |
| 5 | $-1.63$ | $2 \cdot 19$ | -0.79 | -0.74 | 14 | -0.24 | 1.94 | $-0.54$ | $-0.46$ |
| 6 | $-2.38$ | $2 \cdot 89$ | 0.21 | 0.63 | 15 | 0.32 | $4 \cdot 31$ | $2 \cdot 46$ | $2 \cdot 30$ |
| 7 | $-2.38$ | $4 \cdot 34$ | $2 \cdot 19$ | $2 \cdot 31$ | 16 | -0.41 | $5 \cdot 05$ | $3 \cdot 25$ | $3 \cdot 29$ |
| 8 | $-1.90$ | $4 \cdot 91$ | $3 \cdot 44$ | $3 \cdot 31$ | 17 | -0.22 | $3 \cdot 46$ | $1 \cdot 45$ | 1.57 |
| 9 | $0 \cdot 12$ | $5 \cdot 82$ | $4 \cdot 30$ | $4 \cdot 16$ | 18 | $-1.75$ | $3 \cdot 58$ | 1.25 | $1 \cdot 46$ |

* Cf. Fig. 1 and footnote, p. 3303.
significantly affected by these ten additional terms, the molecular outline is altered in the regions where hydrogen peaks may be expected. It is hoped later to refine this structure by a low-temperature technique whose usefulness as regards both accurate bond lengths and location of hydrogen atoms has been demonstrated in the structure of $\alpha$-phenazine (Hirshfeld and Schmidt, Acta Cryst., 1954, 7, 129). In the meantime the general features of the crystal structure and of the molecular shape are quite clear on the basis of the present results.

Fig. 6. The projection of the unit-cell contents on (100).


Fig. 7. The projection of the unit-cell contents on (010). In the lower half of the projection two of the overlapping molecules have been omitted for greater clarity.


Fig. 8. The molecular arrangement in one stack along [001] seen in the (100) projection, showing the short approach distances between neighbouring molecules.


Despite the non-planarity of $3: 4$-benzophenanthrene, the packing arrangement of the molecule in the cell (Figs. 5-8) is similar to that observed in normal aromatic hydrocarbons such as coronene (Robertson and White, $J$., 1945, 607). The molecules are arranged in stacks along the short axis of the unit cell; within each stack the molecules are parallel and inclined to this axis at about $40^{\circ}$, while the mean planes of molecules in different stacks, related by twofold screw axes, are nearly perpendicular.

A very similar arrangement has been postulated for $3: 4-5: 6$-dibenzocarbazole (Iball, Z. Krist., 1936, 95, $A, 282$ ), an overcrowded molecule of the same general shape as $3: 4$ benzophenanthrene: viz., $a=14 \cdot 07, b=6 \cdot 10, c=15.36 \AA$. Space-group, $A 2_{1} 22$; molecular symmetry, 2. There the situation is complicated by the appearance of extra reflections that indicate a superlattice with one axial length four times that of the main lattice. If these are ignored the packing of the molecules as deduced by Iball is seen to be identical with the arrangement, in space-group $B 22_{1}{ }_{2}$, to which $3: 4$-benzophenanthrene approximates. The similarity of the two structures is further emphasised by a general correspondence in the relative intensities of the low-angle spectra, allowance being made for the cyclic permutation of the crystallographic axes suggested by a comparison of the molecular orientations and axial lengths in the two crystals.

It is interesting that the packing in 3:4-benzophenanthrene, while approaching closely to the symmetry of $B 22_{1} 2$, nevertheless deviates slightly from the arrangement of that more symmetrical space-group. Instead of an exact twofold axis parallel to [100] at $v=$ $w=\frac{1}{4}$, the molecule has an approximate twofold axis inclined at about $7^{\circ}$ to [100], with the geometric centre of the molecule located at $u=0.007, v=0.243, w=0.253$. This


Fig. 9. The molecular model viewed at about $30^{\circ}$ to the twofold axis.
would appear to indicate that the lattice energy must be critically dependent on the precise disposition of the molecules in the cell; such a situation might be somewhat surprising were it not for certain short intermolecular contacts that also imply the presence of crystal forces unusual in fully aromatic hydrocarbons.

Within each stack along [001], several pairs of atoms in successive molecules approach to within less than $3.7 \AA$ (Table 2). The shortest intermolecular distances are between

Table 2. Some short contact distances (in $\AA$ ) between molecules separated by unit translation along the c axis.

| Atoms | 8-5' | 8-6' | 16-5' | 7-5' | 16-14' |
| :---: | :---: | :---: | :---: | :---: | :---: |
| Distance | 3•13 | 3.34 | $3 \cdot 48$ | 3.55 | 3•70 |

atoms 8 and $5^{\prime}$, which are $3 \cdot 13 \AA$ apart, and between atoms 8 and $6^{\prime}, 3 \cdot 34 \AA$ apart. These close contacts are apparently related to an unusual valency structure of atoms 5 and 8 , which will be discussed further below. The shortest contacts in the [100] direction are greater than those along [010] (Fig. 5) ; this probably indicates weaker binding forces in the [100] direction, in accordance with the observed cleavage of the crystals parallel to $\{100\}$.

At the present stage of crystallographic refinement the following conclusions about the molecular shape are warranted. The centre-to-centre distance between the overcrowded carbon atoms 1 and 12 is $3.0 \AA ; 3: 4$-benzophenanthrene thus falls into line with previously-studied overcrowded compounds (see Part I, loc. cit.) as regards the minimum approach distance between non-bonded carbon atoms. The manner in which the molecule is distorted from planarity so as to achieve this separation between carbon atoms 1 and 12 may be seen in Fig. 9, which is a perspective view of a model based on the experimental co-ordinates. It is seen that rings $A$ and $A^{\prime}$ are bent in opposite directions out of the mean plane of rings $B$ and $B^{\prime}$.

While the experimental co-ordinates do not conform precisely to the chemicallyprobable molecular symmetry 2 , a molecule having exact twofold symmetry may be
visualised, whose atomic co-ordinates, listed in Table 3, deviate from the experimental values by a root-mean-square difference of only $0.03 \AA$. This deviation is probably not significant, since structure factors calculated for the symmetrical model lead to $R$ factors for the $h k 0,0 k l$, and $h 0 l$ zones of $0 \cdot 23,0 \cdot 25$, and $0 \cdot 27$, which are not much higher than for

Table 3. Atomic co-ordinates (in $\AA_{\text {) }}$ ) adjusted to conform with molecular symmetry 2.

| Atom | $x$ | $y$ | $z$ | Atom | $x$ | $y$ | $z$ | Atom | $x$ | $y$ | $z$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 1 | $1 \cdot 76$ | 2.05 | 0.85 | 7 | $-2.43$ | $4 \cdot 34$ | $2 \cdot 19$ | 13 | 0.40 | $2 \cdot 54$ | 0.54 |
| 2 | $2 \cdot 32$ | $1 \cdot 18$ | 0.03 | 8 | $-1.88$ | $4 \cdot 93$ | $3 \cdot 38$ | 14 | $-0.24$ | 1.92 | $-0.47$ |
| 3 | $1 \cdot 76$ | 0.68 | $-1.23$ | 9 | 0.18 | $5 \cdot 80$ | $4 \cdot 31$ | 15 | 0.33 | $4 \cdot 32$ | $2 \cdot 42$ |
| 4 | 0.44 | $1 \cdot 07$ | $-1.36$ | 10 | $1 \cdot 55$ | 6.01 | $4 \cdot 46$ | 16 | $-0.42$ | $5 \cdot 03$ | $3 \cdot 30$ |
| 5 | $-1.63$ | 2.21 | $-0.84$ | 11 | $2 \cdot 27$ | $5 \cdot 42$ | $3 \cdot 32$ | 17 | $-0.22$ | $3 \cdot 47$ | $1 \cdot 42$ |
| 6 | $-2 \cdot 33$ | $2 \cdot 88$ | $0 \cdot 22$ | 12 | $1 \cdot 77$ | $4 \cdot 62$ | $2 \cdot 39$ | 18 | $-1.76$ | $3 \cdot 57$ | $1 \cdot 27$ |

the experimental co-ordinates. Accordingly, in discussing the shape of the molecule and the manner in which it is deformed, we shall refer exclusively to the dimensions of the symmetrical model. On the assumption that the molecule does in fact possess symmetry 2 , the observed deviations from this symmetry provide an estimate of the accuracy of the present analysis. Thus the experimental co-ordinates may be regarded as subject to a

Fig. 10. Dimensions of the symmetrical model fitted to the experimental co-ordinates. The bond lengths (in $\AA$ ) and the bond angles are indicated on the right, the theoretical mobile bond orders (Berthier et al., loc. cit.) on the left. The circled figures are the distances of the atoms from a reference plane through atoms 13, 15, 17, 18.

root-mean-square error of $0.05 \AA$, which is then also the root-mean-square error of the averaged bond lengths (except for bond 17-18, which occurs only once) shown in Fig. 10.

The most striking feature of this structure is the non-planarity of the " benzene" rings. Atoms 1, 2, 3, 4, 14, 13, constituting ring $A$, deviate from the best plane through these atoms by a root-mean-square distance of over $0.05 \AA$ and a maximum of $0.08 \AA$, the deviations being systematic so that the ring may be regarded as folded, through an angle of $10^{\circ}$, along the diagonal joining atoms 3 and 13 . The distortion of ring $B$ is even greater, its atoms deviating from their best plane by a root-mean-square distance of $0.08 \AA$ and a maximum of $0 \cdot 10 \AA$. Here again the principal mode of deformation consists in a folding of the ring, through an angle of $14^{\circ}$ along the diagonal 5-17. It seems certain that this folding of rings $A$ and $B$, together with the corresponding folding of the symmetry-related rings $A^{\prime}$ and $B^{\prime}$, accounts for most of the distortion of the molecular skeleton by which relief from overcrowding is achieved. Any other deformations of the skeleton appear to be of secondary importance, atoms $1,2,3,4$ being coplanar to within a root-mean-square deviation of $0.02 \AA$, while the groups comprising atoms $3,4,14,5,17,13$ and atoms 5,6 , $18,7,8,17$ are each planar to within a root-mean-square deviation of $0.04 \AA$.

The molecular properties likely to be affected by the distortions of the aromatic system include resonance energy, dipole moment, valency structure, and chemical reactivity; these will be discussed in turn.

The results of calculations, by the method of molecular orbitals, of theoretical resonance energies of all four-ring aromatic hydrocarbons $\mathrm{C}_{18} \mathrm{H}_{12}$ (Berthier, Coulson, Greenwood, and A. Pullman, Compt. rend., 1948, 226, 1906) are quoted in Table 4. The computations, for planar molecules, were made both with and without allowance for overlap between
adjacent atomic orbitals, the results being given in terms of the empirical constants $\gamma$ and $\beta$ respectively. The corresponding experimental values have been determined by Magnus and Becker (Erdöl und Kohle, 1951, 4, 115). By combining the theoretical and experimental results for all these compounds except $3: 4$-benzophenanthrene, we obtain values for $\gamma$ and $\beta$ of $37.64 \pm 1.0 \mathrm{kcal} . /$ mole and $18.61 \pm 0.15 \mathrm{kcal} . /$ mole, respectively. The resonance energy per mole of the planar model of 3:4-benzophenanthrene is thus

Table 4. Comparison of experimental and theoretical resonance energies of all

| Compound | four-ring aromatic hydrocarbons $\mathrm{C}_{18} \mathrm{H}_{12}$. |  |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | Exptl. resonance energy (kcal./ mole) | Quantumwith overlap | echanical energy: without overlap ( $\beta$ ) | $\begin{gathered} \gamma^{+} \\ (\mathrm{kcal} . / \\ \text { mole) } \end{gathered}$ | $\begin{gathered} \beta+\ddagger \\ (\text { kcal. } \\ \text { mole) } \end{gathered}$ |  | ical $\stackrel{\text { energy }}{=18.61}$ |
| 3:4-Benzophenanthrene | 127.7 | 3.5840 | 7-1874 | $35 \cdot 63$ | 17.77 | 134.9 | 133.8 |
| Tetracene | 129.6 | $3 \cdot 3400$ | 6.9316 | 38.80 | 18.70 | 125.7 | 129.0 |
| 1: 2-Benzanthracene | $130 \cdot 4$ | 3.5027 | $7 \cdot 1012$ | 37.23 | 18.36 | $131 \cdot 8$ | 132.2 |
| Chrysene | $134 \cdot 4$ | 3.5872 | $7 \cdot 1900$ | 37.47 | 18.69 | $135 \cdot 0$ | 133.8 |
| Triphenylene | $135 \cdot 8$ | $3 \cdot 6658$ | 7.2745 | 37.05 | 18.67 | 138.0 | $135 \cdot 4$ |

* Magnus and Becker, loc. cit.
$\dagger$ Berthier et al., loc. cit.
$\ddagger$ Values derived from comparison of corresponding experimental and theoretical resonance energies.
calculated as $134.9 \pm 3.5 \mathrm{kcal}$., from $\gamma$, or $133.8 \pm 1 \cdot 1 \mathrm{kcal}$., from $\beta$; both values are greater than the experimental value of 127.7 kcal . The energy defect of 6 or 7 kcal . must be related to the energy of intramolecular repulsion, neglected in the theoretical calculations, and to the deformation of the hypothetically-planar ring system.

The dipole moment of $3: 4$-benzophenanthrene has been reported as 0.7 D (Bergmann, Fischer, and B. Pullman, J. Chim. phys., 1951, 48, 356). The major contribution to this moment probably comes from the dipoles of the $\mathrm{C}-\mathrm{H}$ bonds, which do not cancel in the deformed structure as they would in a regular planar molecule.

The folding of the rings is reflected in the abnormal bond angles at the carbon atoms on the axes of fold; the bond angles at the secondary atoms 3 and 5 are $110^{\circ}$ and $111^{\circ}$, while the corresponding angles at the tertiary carbon atoms 13 and 17 are $116^{\circ}$ and $114^{\circ}$. Although the individual values may not be too reliable, it is significant that they all deviate in the same direction from the trigonal bond angle of $120^{\circ}$ and that those at the secondary carbon atoms approach the tetrahedral angle of $109 \frac{1}{2}^{\circ}$. Thus it appears that the valency electrons of these secondary atoms do not occupy pure $\left(s p^{2}\right) p$ orbitals, appropriate to a planar aromatic structure, but more nearly approximate to the tetrahedral $s p^{3}$ hybridisation. This must imply an enhanced localisation of the $\pi$ electrons that these atoms contribute to the aromatic system and may be expressed in a greater polarisability than would be expected in an undeformed molecule. Whether or not this localisation involves a decoupling of the electron spins should be deducible from magnetic data, which are unfortunately not available. A knowledge of the hydrogen positions would be of considerable value in specifying more exactly the type of hybridisation present.

The increased free valence that may be expected to accompany the apparent shift towards tetrahedral hybridisation probably accounts for the short approach distance between atoms 5 and 8 of neighbouring molecules in the crystal, which may indicate a tendency towards additional bond formation on the part of these " unsaturated " atoms. On the basis of the co-ordinates of the symmetrical model it is found that the line joining atoms 5 and 8 of neighbouring molecules makes, at one end, angles of $85^{\circ}$ and $112^{\circ}$ with bonds 5-6 and 5-14 respectively and, at the other end, angles of $94^{\circ}$ and $91^{\circ}$ with bonds $8-7$ and 8-16 respectively.

The known chemical properties of $3: 4$-benzophenanthrene support the above view of the valency structure since the increased availability of negative charge at the distorted carbon atoms provides an explanation for the preferential substitution at the 5 -position by bromine, nitric acid, and the Friedel-Crafts reagent (Newman and Kosak, J. Org. Chem., 1949, 14, 375). Also the tetrahedral hybridisation already evident at the reactive carbon
atom in the unperturbed molecule would favour the formation of the activated complex, in which, supposedly, the entering substituent and the outgoing hydrogen are simultaneously bound to this carbon atom in a tetrahedral arrangement (e.g., Waters, J., 1948, 727). We

Table 5. List of observed and calculated structure factors.
(a) $\mathbf{F}(h k 0)$ referred to origin $\left(\frac{1}{1} 00\right)$.


* Signs determined on $X$-RAC by means of the " negative-background criterion."
$\dagger$ Presumed secondary-extinction error.

Table 5. (Continued.)

| (b) $\mathbf{F}(0 k l)$ referred to origin |  |  |  |  |  |  |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 0kl | $\mathrm{F}_{0}$ | $\mathrm{F}_{\mathrm{C}}$ | 0 kl | $\mathrm{F}_{0}$ | $\mathrm{Fc}_{\mathrm{c}}$ | 0kl | $\mathrm{F}_{0}$ | $\mathrm{F}_{\mathrm{C}}$ | 0kl | Fo | $\mathrm{F}_{\mathrm{c}}$ |
| 002 | 13 | -11 | 012 | 8 | $-10$ | 073 | 0 | 1 | 015 | 10 | 15 |
| 4 | 0 | 3 | 2 | 20 | 20 | 8 | 6 | -10 | 2 | 3 | 7 |
| 6 | 12 | $-20$ | 3 | 11 | 10 | 9 | 18 | -16 | 3 | 0 | $-4$ |
|  |  |  | 4 | 23 | -24 | 10 | 8 | 5 | 4 | 0 | 1 |
|  |  |  | 5 | 35 | 34 | 11 | 11 | - 9 | 5 | 8 | 6 |
| 011 | 0 | - 2 | 6 | 43 | -40 | 12 | 8 | 5 |  |  |  |
| 2 | 10 | 13 | 7 | 26 | 26 |  |  |  | 016 | 0 | - 2 |
| 3 | 8 | -11 | 8 | 9 | $-5$ |  |  |  | 2 | 3 | - 6 |
| 4 | 14 | $-5$ | 9 | 3 | - 3 | 014 | 0 | - 2 | 3 | 0 | 1 |
| 5 | 25 | 18 | 10 | 8 | -13 | 2 | 12 | -15 | 4 | 3 | 4 |
| 6 | 27 | -23 | 11 | 18 | -18 | 3 | 5 | $-2$ | 5 | 0 | $-1$ |
| 7 | 6 | 7 | 12 | 9 | -12 | 4 | 15 | 17 | 6 | 0 | $-1$ |
| 8 | 9 | 11 | 13 | 6 | $-9$ | 5 | 14 | 19 | 7 | 0 | 1 |
| 9 | 0 | $-1$ |  |  |  | 6 | 30 | 37 | 8 | 0 | 1 |
| 10 | 4 | 1 |  |  |  | 7 | 14 | 15 | 9 | 4 | 2 |
| 11 | 8 | 6 | 013 | 10 | -11 | 8 | 0 | $-2$ |  |  |  |
| 12 | 0 | 1 | 2 | 2 | 4 | 9 | 9 | $-4$ | 017 | 0 | 5 |
| 13 | 0 | $-1$ | 3 | 7 | 6 | 10 | 0 | - 4 | 2 | 0 | 2 |
| 14 | 8 | 8 | 4 | 10 | - 7 | 11 | 8 | - 8 | 3 | 0 | - 2 |
| 15 | 5 | $-5$ | 5 | 10 | 8 | 12 | 11 | 7 | 4 | 0 | 2 |
|  |  |  | 6 | 8 | $-5$ | 13 | 7 | - 6 | 5 | 5 | -10 |

(c) $\mathrm{F}(h 0 l)$ referred to origin ( $00 \frac{1}{4}$ ).

| $h 0 l$ | $\mathrm{F}_{0}$ | Fc | h0l | $\mathrm{F}_{0}$ | $\mathrm{F}_{\mathrm{c}}$ | $h 0 l$ | $\mathrm{F}_{0}$ | $\mathrm{F}_{\mathrm{c}}$ | h0l | Fo | $\mathrm{F}_{\mathrm{c}}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 101 | 24 | 24 | 402 | 7 | 2 | 803 | 0 | 3 | 105 | 3 | 7 |
| 2 | 3 | - 2 | 5 | 0 | 1 | 9 | 20 | 20 | 2 | 5 | $-4$ |
| 3 | 8 | -13 | 6 | 16 | 19 | 10 | 0 | - 2 | 3 | 6 | $-7$ |
| 4 | 0 | - 2 | 7 | 0 | $-5$ | 11 | 25 | -28 |  |  |  |
| 5 | 7 | $-9$ | 8 | 11 | 16 | 12 | 0 | $-1$ |  |  |  |
| 6 | 8 | - 7 | 9 | 5 | $-7$ | 13 | 6 | 6 |  | 11 | -10 |
| 7 | 4 | -4 | 10 | 15 | -16 |  |  |  |  | 11 | $-10$ |
| 8 | 6 | 6 | 11 | 6 | 9 | 104 | 0 | $-5$ | 3 4 | 0 9 | 3 5 |
| 9 | 0 | 2 |  |  |  | 2 | 0 | $-5$ | 4 | 9 | 5 |
| 10 | 0 | 1 | 103 | 19 | 19 | 3 | 6 | $-7$ |  |  |  |
| 11 | 5 | 1 | 2 | 2 | -9 | 4 | 0 | 7 | 506 | 0 | 2 |
|  |  |  | 3 | 46 | -42 | 5 | 0 | 0 | 6 | 3 | $-5$ |
|  |  |  | 4 | 4 | $-2$ | 6 | 0 |  | 7 | 0 | $-1$ |
| 102 | 4 | - 2 | 5 | 28 | 19 | 7 | 0 | 2 | 8 | 3 | 6 |
| 2 | 36 | -27 | 6 | 0 | 4 | 8 | 0 | -1 |  |  |  |
| 3 | 3 | 9 | 7 | 0 | 3 | 9 | 8 | 5 |  |  |  |

prefer this explanation to the one based on an analogy with the reactivity of phenanthrene at the 9 -position, which would in any case not account for the preferential attack on the 5 - rather than the 6 -position in 3:4-benzophenanthrene. Two disubstitution products were also found in the nitration reaction but their constitution was not investigated further. Unless primary substitution fundamentally changes the structure of the molecule we should expect, on the above arguments, that the second substituent has entered at the 8 -, $3-$, or 10 -position.

Although the present bond-length measurements are not sufficiently accurate to warrant detailed analysis, their general tendency is in accord with the proposed interpretation of the molecular valency structure. Some idea of the bond lengths to be expected on this hypothesis may be derived from a consideration of the eight biradical structures (Fig. 11) suggested by the tetrahedral bond angles at atoms $3,5,8$, and 10 . Although there is no evidence that the molecule is, in fact, a biradical, we should expect similar bond-length variations to occur even if the electron spins are not uncoupled since there are one Dewar long-bond and two ionic structures corresponding, in the distribution of single and double bonds, with each of the proposed biradical structures and these would probably contribute strongly, along with the eight Kekulé-type structures, to the resultant bond structure of a molecule with no unpaired electrons. Thus, the length of the central 17-18 bond, which is certainly exaggerated by the present data, is partly explained by the fact that it appears single in all the suggested non-Kekulé formulæ, as does bond $5-14$, whose measured length
is $1.47 \AA$. These formulæ would also indicate three-fourths double-bond character for bond 1-2, measured as $1.32 \AA$ long. It may also be significant that bond $5-6$, with a measured length of $1 \cdot 44 \AA$ despite its theoretical mobile-bond order of 0.762 calculated by Berthier et al. (loc. cit.), for a planar molecule, appears as a double bond in just half the biradical, or pseudo-biradical, formulæ under consideration. Thus the measured bond lengths, insofar as they are reliable, seem generally consistent with the other evidence as to the type of electronic structure to be associated with the distorted benzophenanthrene molecule.

The conclusions reached here provide a coherent interpretation of the known physical and chemical properties of 3:4-benzophenanthrene, as well as indicating several directions in which further work is needed. On the crystallographic side, comparison of these results with the analysis of the crystal structure of $5: 4^{\prime}$-dimethyl-3:4benzophenanthrene, which is already well under way, will be of particular interest.

The crystallographic evidence indicates resolution of 3:4-benzophenanthrene into $(+)$ - and

Fig. 11. Four of the eight biradical formula suggested by the tetrahedral bond configurations. The remaining four are related to these by the operation of the molecular twofold axis of symmetry.





 $(-)$-crystals. We have not yet succeeded in growing crystals either of large enough size or with sufficiently well-developed $h k l$ faces to take advantage of this spontaneous resolution for the measurement of the optical activity of 3:4-benzophenanthrene; experiments in this direction are in progress.

Experimental.-Good crystals of 3:4-benzophenanthrene can be obtained by slow cooling of a solution in ethanol. They are colourless needles, elongated along [001], whose most prominent faces are $\{100\}$ and $\{010\}$, with $\{110\}$ and $\{011\}$ frequently present. The crystals show a cleavage plane parallel to $\{100\}$.

All zonal reflections obtainable with $\mathrm{Cu}-\mathrm{K} \alpha$ radiation were recorded on Weissenberg photographs about the three principal axes, and their intensities measured visually by comparison with a series of standard spots. The roughly cylindrical outline of the crystals made spot-shape and absorption corrections in the $h k 0$ zone unnecessary; for photographs about the [100] and [010] axes, a crystal was cut to approximately cubic shape. The absolute scale, originally estimated by Wilson's statistical method (Nature, 1942, 150, 151), was later found, on comparison with the scale of the final calculated structure factors, to be about $10 \%$ higher in $\mathrm{F}(h k l)$ that the scale originally adopted. The accuracy of the observed structure factors is thought to be of the order of $10 \%$.

We are grateful to Professor M. S. Newman and to Dr. J. Szmuskowicz for gifts of crystals of 3:4-benzophenanthrene, and to Mr. F. L. Hirshfeld for many helpful suggestions during the preparation of this paper. One of us (G. M. J. S.) is greatly indebted to Professor R. Pepinsky for the use of $X-\mathrm{RAC}$, and to the Rockefeller Foundation for financial support, during his stay at State College, Pa. We also acknowledge the help of Mr. B. Smilovits with the calculations and of Mr. R. Cohen with the preparation of the diagrams. This work was carried out during the tenure of the Porter scholarship of the University of Cape Town.

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[^0]:    * Part II, preceding paper.
    $\dagger$ In this paper the numbers used in references to the carbon atoms are those shown in Fig. 1 and not those strictly related to the name 3:4-benzophenanthrene.

